Study of the discovery potential of next generation LSW experiments

Paola Arias, DESY 6th PATRAS Workshop on Axions, WIMPS and WISPS 2010

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Outline

- ➡ WISPs intro.
- Current bounds, (Axion-light particles)
- Laboratory experiments: Photonregeneration experiments (LSW).
- Next generation of experiments.
- Conclusions.



Weakly Interacting Sub-eV Particles

- WISPs appear naturally in many extensions of the Standard Model.
- Hidden sectors: light particles weakly coupled with the visible sector. Among others: AXIONS¹ HIDDEN U(1)'s MINI-CHARGED

$$\mathcal{L}_{int} = g_{a\gamma\gamma} a F_{\mu
u} ilde{F}^{\mu
u}$$

$$\mathcal{L}_{mix} = \frac{1}{2} \chi F_{\mu\nu} B^{\mu\nu}$$

$$Q_{MCP} = \frac{e_h}{e} \chi$$

Nice introduction from I. Antoniadis, J. Jaeckel



1 Peccei- Quinn (1977), Wilczek (1978), Weinberg (1977) 2 Holdom (1985)

Where to search?

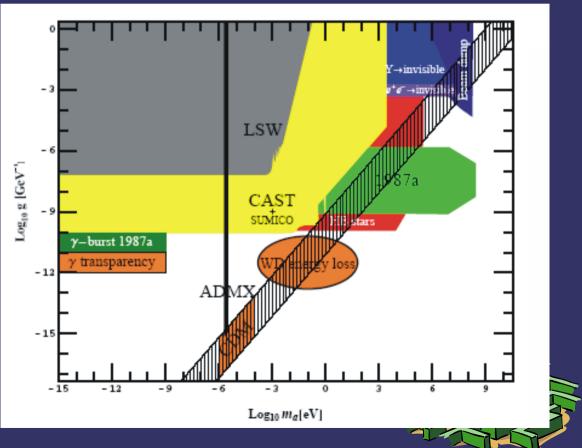
Fortunately, our universe is big factory of WISPs: Several Astrophysical experiments constraint axion-like particles (ALPs)

Hints from puzzling Astrophysical observations:

• Cosmic photon oscillation Into ALPs in random magnetic fields in galaxy Clusters $m_0 \le 10^{-12}$ eV. ¹ ?

• Transparency of the Universe due high energy $\gamma \rightarrow \phi$ conversion². (talk from M. Roncadelli)

¹ Burrage et al (2009)
 ² A. De Angelis et al (2007)

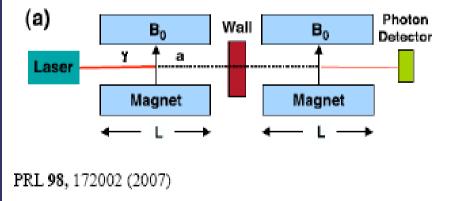


Boom of Laboratory Experiments

2005 - 2007 PVLAS OBSERVATION – RETRACTION of vacuum birefringence and dichroism in laser polarization experiments.

Prompt a boom of theoretical research (WISPs) and more importantly: optical precision experiments, in particular photon regeneration experiments

Via Primakoff effect we expect to detect regenerated photons after the wall



Sensitive:



Light Shining through a wall

Advantages: Sensible to ALPs, Hidden U(1)'s and MCPs, relatively simple. Disadvantage: Lose of coherence for higher masses.

At first order in (g B L):

$$P_{\gamma \to a} = \frac{(g_{a\gamma\gamma} B \cdot L \cdot \omega)^2}{|k_a|^2} |F(q)|^2 \qquad q = |k_{\gamma} - k_a| \sim \frac{m^2}{2\omega}$$
Strong B xL
dependence!
$$F(q) = \frac{1}{L} \int_0^L e^{iqz} dz = \frac{2}{qL} \sin\left(\frac{qL}{2}\right) \qquad \text{Raffelt \& Stodolsky (1988)} \\ \text{Van Bibber et al (1987)}$$
Probability peaks for F(q) $\rightarrow 1$ $\qquad q \to 0$
Maximal sensibility for small masses.
LSW bounds $m_{\phi} < 10^{-3} \text{ eV}, \qquad g_{\phi\gamma\gamma} \lesssim 7 \times 10^{-8}$ ALPS Collaboration

Next Generation of LSW experiments

Technological and Theoretical improvements are needed!

Resonant Enhanced FP Cavities: Probability of conversion enhanced by $\frac{2}{\pi^2} \frac{\mathcal{F}_{\gamma} \mathcal{F}_{a}}{\pi^2}$, \mathcal{F} : Finesse of the cavity

Actual Finesse achivable 10⁴ - 10⁵

Powerful laser and detector technology

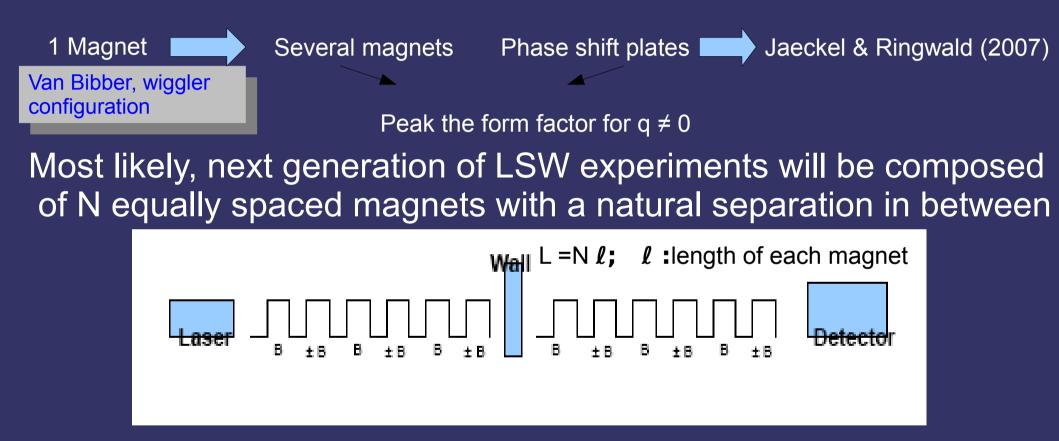
150 kW laser power **Detector efficiency**

Reciclyng of superconducting dipole magnets (HERA, TEVATRON, LHC). Improve B x L!!!!!!



2) Hoogeveen & Ziegenhagen, Nucl. Phys B 358 (1991); Sikivie, Tanner & van 98 (2007)

Second Generation of LSW experiments



We can decompose the form factor:

$$\int_0^L dx' e^{iqx'} = \int_0^l dx' e^{iqx'} + \int_{l+\Delta}^{2l+\Delta} dx' e^{iqx'} + \dots + \int_{(N-1)(l+\Delta)}^{l+(N-1)(l+\Delta)} dx' e^{iqx'},$$



For the normal non-alternating configuration Works like we were adding a phase of $q\Delta$ for each magnet

$$F(q) = \int_0^\ell dx \ e^{iqx} + \int_\ell^{2\ell} dx \ e^{iq(x+\Delta)} + \ldots + \int_{(N-1)\ell}^{N\ell} dx \ e^{iq(x+(N-1)\Delta)}$$

$$F(q) = \sum_{s=0}^{N-1} \int_{s\ell}^{(s+1)\ell} dx \, \exp\left[iq \left(x + s\Delta\right)\right]$$

See : "Extending the reach of axion-photon regeneration experiments towards larger masses with PSP" Jaeckel & Ringwald, PLB(2007).

As with PSP, we are able to scan a wide region of ALPs masses verying the gaps between the magnets



Without introducing extra losses!

As a result, we obtain modified form factors:

Normal N-magnets configuration:

$$F(q) = \frac{2}{qL} \sin\left(\frac{qL}{2N}\right) \frac{\sin\left(\frac{qN}{2}\left(\frac{L}{N} + \Delta\right)\right)}{\sin\left(\frac{q}{2}\left(\frac{L}{N} + \Delta\right)\right)}.$$

Alternating magnets of n subgroups:

$$F(q) = \begin{cases} \frac{2q}{L} \sin\left(\frac{qL}{2N}\right) \frac{\sin\left(\frac{qN}{2}(L/N+\Delta)\right)}{\sin\left(\frac{q}{2}(L/N+\Delta)\right)} \tan\left(\frac{qN}{2n}\left(\frac{L}{N}+\Delta\right)\right), & n \text{ even} \\ \frac{2q}{L} \sin\left(\frac{qL}{2N}\right) \frac{\cos\left(\frac{qN}{2}(L/N+\Delta)\right)}{\sin\left(\frac{q}{2}(L/N+\Delta)\right)} \tan\left(\frac{qN}{2n}\left(\frac{L}{N}+\Delta\right)\right), & n \text{ odd} \end{cases}$$

Maximizing the form factor, we can improve the sensibility for a certain mass range:

Example, Normal

$$\frac{qL}{2N}\left(1+\frac{N\Delta}{L}\right) = k\pi,$$



Prospects for ALPs

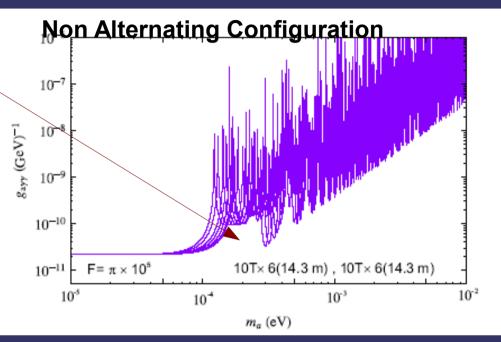
2.8

6+6 configuration LHC

Res. Cavities, Laser Power Detector

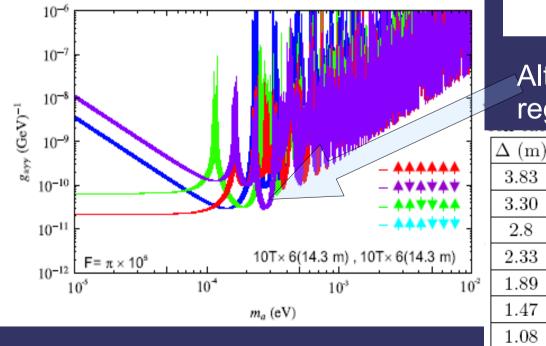
$$g^{min}_{\phi\gamma\gamma} \lesssim 3 \times 10^{-11}$$

	$m_{\phi} \times 10^{-4} \text{ eV}$	Δ (m)
	3	$\left(15.4 \right)$
	3.2	11.8
	3.5	7.5
	3.7	5.23
	3.9	3.3
	4.1	1.6



Alternating configurations help to fill the region of small masses





Conclusions

- We have derived the corrected form factors for the next generation of LSW experiments.
- According to the goals of the individual experiment, an appropriate gap between the magnets should be selected, or at least taken into account for realistic predictions.
- Restrictions: Cavity size, aperture of the magnets. Available magnets allow cavities of the order of 100-150 m.
- New experimental ideas to "exploit" the gaps?
 Also in MCPs, Chameleons?

